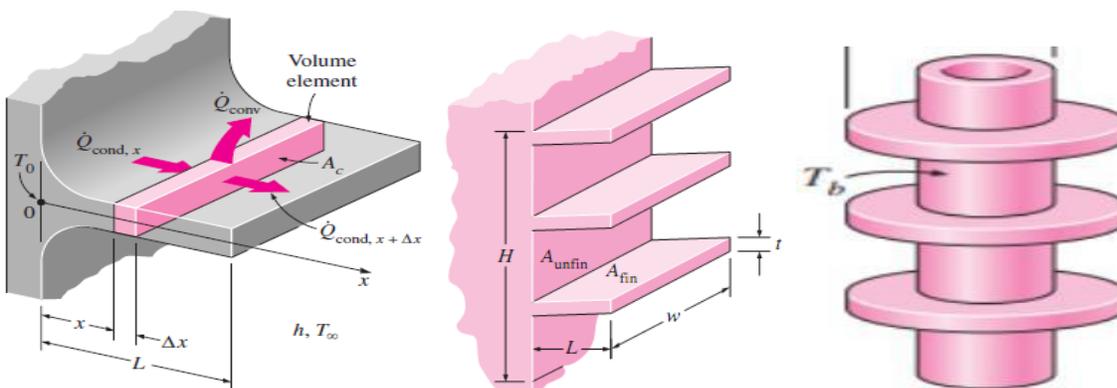




Fin Equation:

Consider a volume element of a fin at location x having a length of Δx , cross-sectional area of A_c , and a perimeter of p . Under steady conditions, the energy balance on this volume element can be

$$\left(\begin{array}{l} \text{Rate of heat} \\ \text{conduction into} \\ \text{the element at } x \end{array} \right) = \left(\begin{array}{l} \text{Rate of heat} \\ \text{conduction from the} \\ \text{element at } x + \Delta x \end{array} \right) + \left(\begin{array}{l} \text{Rate of heat} \\ \text{convection from} \\ \text{the element} \end{array} \right)$$



Or

$$\dot{Q}_{\text{cond},x} = \dot{Q}_{\text{cond},x+\Delta x} + \dot{Q}_{\text{conv}}$$

Where

$$\dot{Q}_{\text{conv}} = h(p \Delta x)(T - T_{\infty})$$

Substituting and dividing by Δx , we obtain: [perimeter: $p = 2(w+t)$]

$$\frac{\dot{Q}_{\text{cond},x+\Delta x} - \dot{Q}_{\text{cond},x}}{\Delta x} + hp(T - T_{\infty}) = 0$$



Taking the limit as $\Delta x \rightarrow 0$ gives

$$\frac{d\dot{Q}_{\text{cond}}}{dx} + hp(T - T_{\infty}) = 0$$

In general, the cross-sectional area A_c and the perimeter p of a fin vary with x , which makes this differential equation difficult to solve. In the special case of *constant cross section* and *constant thermal conductivity*, the differential equation reduces to,

$$\frac{d^2\theta}{dx^2} - a^2\theta = 0$$

Where,

$$a^2 = \frac{hp}{kA_s}$$
$$a = \sqrt{\frac{hp}{kA_s}}$$

and ($\theta = T - T_{\infty}$) is the *temperature excess*. At the fin base ($\theta_b = T_b - T_{\infty}$)



$$M = \sqrt{hp k A_c} \theta_b$$

At the fin tip we have several possibilities,

1- Infinitely Long Fin : ($T_{\text{fin tip}} = T_{\infty}$),

at the fin base where $x = 0$ the value of θ will be θ_b . Let, p is the perimeter, and A_c is the cross sectional area. The rates of heat transfer for both cases are given to be

The temperature distribution along the fin **for very long fins is**,

$$\frac{T_x - T_{\infty}}{T_b - T_{\infty}} = e^{-ax} = e^{-x\sqrt{hp/kA_c}}$$

The rates of heat transfer:

$$Q_{\text{long fin}} = -kA_c \frac{dT}{dx} \Big|_{x=0} = \sqrt{hPkA_c} (T_b - T_{\infty})$$



Alternatively, the rate of heat transfer from the fin could also be determined by considering heat transfer from a differential volume element of the fin and integrating it over the entire surface of the fin. That is,

$$\dot{Q}_{\text{fin}} = \int_{A_{\text{fin}}} h[T(x) - T_{\infty}] dA_{\text{fin}} = \int_{A_{\text{fin}}} h\theta(x) dA_{\text{fin}}$$

2- Insulated Fin Tip: ($Q_{\text{fin tip}} = 0$)

for fins with negligible heat transfer at the fin are given by,

- The condition at the fin base remains the same, the relation for **the temperature distribution is,**

$$\frac{T_x - T_{\infty}}{T_b - T_{\infty}} = \frac{\cosh a(L - x)}{\cosh aL}$$

- The rate of heat transfer from the fin can be determined again **from Fourier's law of heat conduction,**

$$Q_{\text{insulated tip}} = -kA_c \frac{dT}{dx} \Big|_{x=0} = \sqrt{hPkA_c} (T_b - T_{\infty}) \tanh aL$$



- ❖ Note that the heat transfer relations for the very long fin and the fin with negligible heat loss at the tip differ by the factor **$\tanh aL$** , which approaches **1 as L becomes very large**.

3- Fin subjected to Convection (and Radiation) at Tip:

Fins exposed to convection at their tips **can be treated as fins with insulated tips** by using the **corrected length (L_c)** instead of the actual fin length.

$$L_c = L + \frac{A_c}{p}$$

Primary References

- Holman, J.P., *Heat Transfer, 10th Edition, McGraw-Hill.*
- Kern, D.Q., *Process Heat Transfer, McGraw-Hill.*

Additional Reference

- Çengel, Y.A. & Ghajar, A.J., *Heat and Mass Transfer: Fundamentals and Applications, McGraw-Hill.*